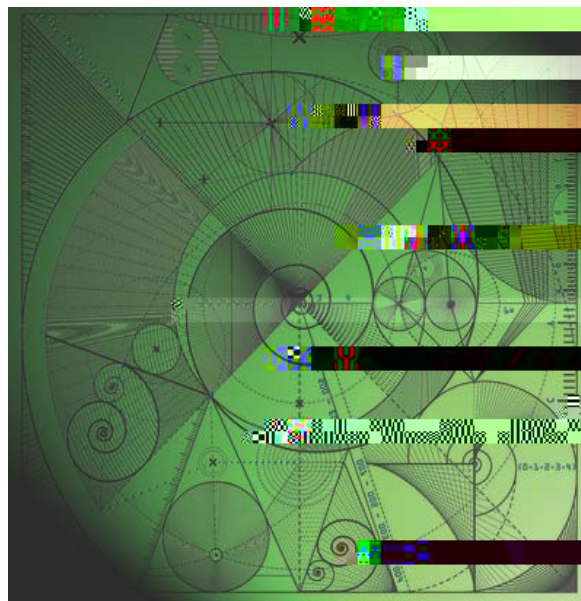


Numerically stable computation of embedding formulae for scattering by polygons

by

A. Gibbs,

and A. Moiola

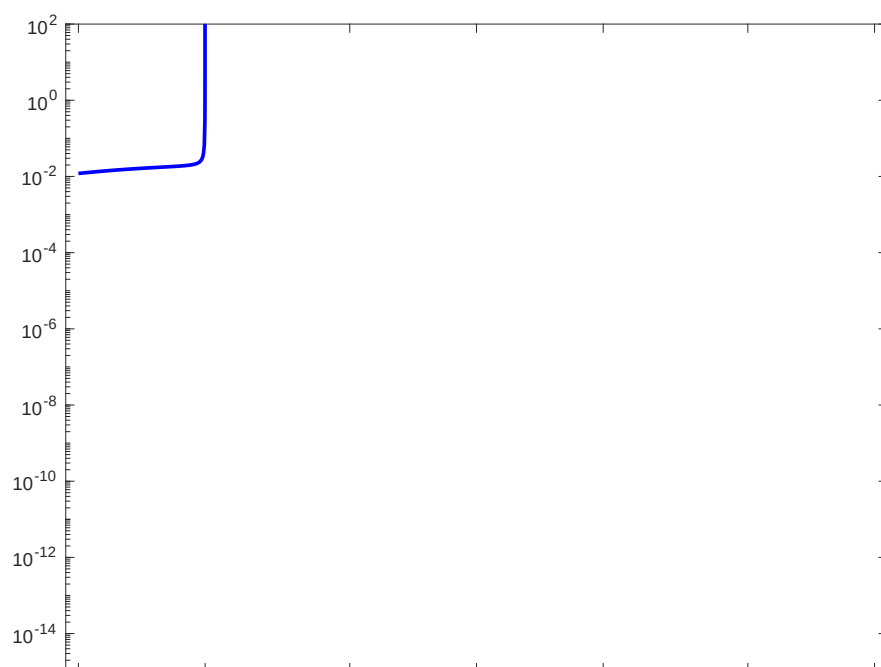


Numerically stable computation of embedding formulae for scattering by polygons

A. Gibbs · S. Langdon · A. Moiola

Abstract For problems of the acoustic scattering by polygonal obstacles embedding formulae provide a useful means of computing the far field coefficient induced by any incident plane wave given the far field coefficient of a relatively small set of canonical problems—the number of such problems to be solved depends only on the geometry of the scatterer. Whilst the formulae themselves are exact in theory any implementation will incur numerical error from the method used to solve the canonical problems—this error can lead to numerical instabilities. Here we present an effective app

nce $|s_n(p_{\theta}/\gamma)|$ is symmetric about the points $x \in \Theta$ it follows that $|s_n(p(\theta + \theta_0)/\gamma)| = |s_n(p(\theta - \theta_0)/\gamma)|$
ence $|s_n(p_{\theta}(\theta + \theta_0))| \leq$



we now have

$$\|D(\theta, \alpha) - \mathcal{E}_{\mathcal{P}} D(\theta, \alpha, \theta, N_T)\| = \left\| \frac{\theta - \theta}{\Lambda(\theta, \alpha)} \sum_{m=1}^M \sum_{n=1}^{N_T} \frac{(\theta - \theta)^{n-1}}{n} B_m(\alpha) \frac{\partial^n \hat{D}}{\partial \theta^n}(\theta, \alpha_m) - b_m(\alpha) \frac{\partial^n}{\partial \theta^n} \mathcal{P} \hat{D}(\theta, \alpha_m) + \mathcal{R}_{\mathcal{P}} D(\theta, \alpha, \theta, N_T) \right\|,$$

where

$$\mathcal{R}_{\mathcal{P}} D(\theta, \alpha, \theta, N_T) = \sum_{m=1}^M B_m(\alpha) \sum_{n=N_T+1}^{\infty} \frac{(\theta - \theta)^{n-1}}{n} \frac{\partial^n \hat{D}}{\partial \theta^n}(\theta, \alpha_m).$$

It follows immediately from Lemma 2 that

$$\left\| \frac{\theta - \theta}{\Lambda(\theta, \alpha)} \right\| \leq \frac{1}{p \|\theta_{\square} - \theta\|}, \quad \text{for } \theta_{\square} \in \Theta_{\square},$$

which compares favourably to (3) when θ is close to $\theta \in \Theta_{\alpha}$ provided that θ is sufficiently far from any $\theta_{\square} \in \Theta_{\square}$ which is enforced by the second condition in (1).

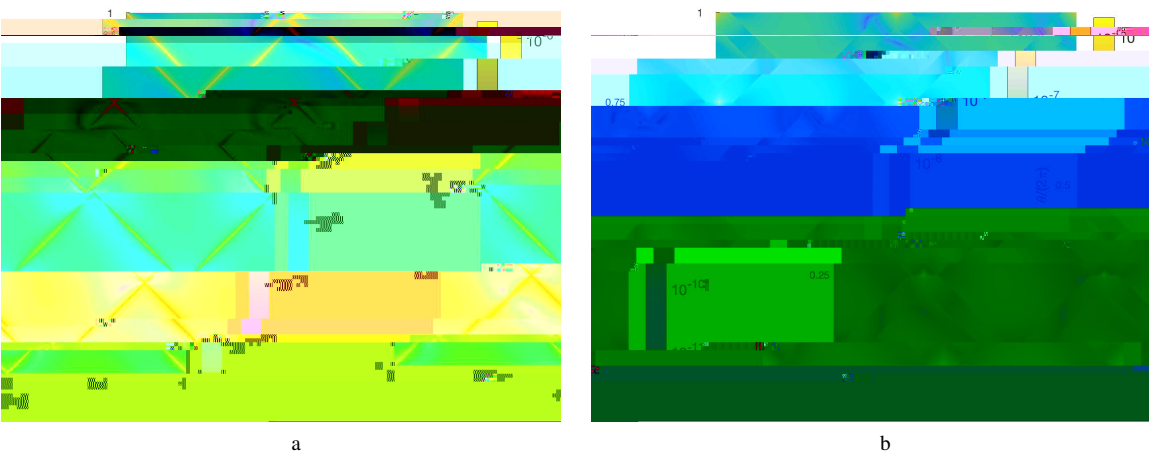
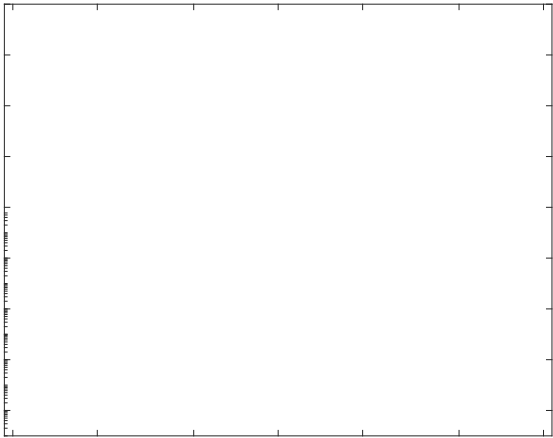


Figure 1: Two plots of relative error for different representations of embedding for feature scattering by a square of side length $k = 10^{-6}$. Plot 'a' depicts accuracy for a naive representation



5 More general incident waves

We now demonstrate how the beddng for uae may be used to approximate the far field coefficient of a far broader class of incident waves than just plane waves by means of a general formula and so the numerical calculation is simpler. In particular we will be interested in incident waves of the following form (see e.g. [1] Definition 5.1) which can be thought of as continuous near combinations of plane waves.

Definition 5.1 (Herglotz wave functions) Given $g_{\text{Herg}} \in L^2(\mathbb{S}^1, \mathbb{C})$ the function

$$u_{\text{Herg}}^i(\mathbf{x}; g_{\text{Herg}}) = \int_{\mathbb{S}^1} g_{\text{Herg}}(\alpha) e^{i\mathbf{x} \cdot \mathbf{d}_\alpha} d\alpha, \quad \text{for } \mathbf{x} \in \mathbb{R}^2,$$

where $\mathbf{d}_\alpha = (\cos \alpha, \sin \alpha)$ is called a *Herglotz wave function* or *Herglotz incident field* with Herglotz kernel $g_{\text{Herg}} \in L^2(\mathbb{S}^1, \mathbb{C})$.

A second concept which we will find useful is the *far-field map* $\mathcal{F}_\infty$ — we first extend the scattering boundary value problem (1.1) to general incident fields. If $u^i \in C^\infty(\mathbb{R}^2)$ is an entire solution of the Helmholtz equation we denote by $u = u^i + u^s$ the solution of the Helmholtz equation (1.1) in the complement of the scatterer with $u = 0$ on $\partial\Omega$ and such that u^s satisfies the outgoing radiation condition — then we write $\mathcal{F}_\infty u^i \in C^\infty(\mathbb{S}^1, \mathbb{C})$ for the far field coefficient of the scattered field u^s . For example in terms of far plane wave notation (1.2) and (1.3) we have $\mathcal{F}_\infty u_\alpha^i = D(\cdot, \alpha)$.

By Definition 5.1 the far field coefficient of a Herglotz wave function depending on Ω can be computed by integrating the plane wave far field coefficient against the Herglotz kernel so it can be approximated using the combined beddng approximation of Definition 5.1.

$$\mathcal{F}_\infty u_{\text{Herg}}^i(\cdot; g_{\text{Herg}})(\theta) = \int_{\mathbb{S}^1} g_{\text{Herg}}(\alpha) D(\theta, \alpha) d\alpha \approx \int_{\mathbb{S}^1} g_{\text{Herg}}(\alpha) \mathcal{E}_{\mathcal{D}}^{\otimes} D(\theta, \alpha; N_T) d\alpha. \quad (5.1)$$

In practice



. The statistics DOFs per side of the scatterer are the same as for the standard errors observed in Figure 1.

w/c can be obtained by solving on M problems of plane wave incidence—Figure — shows the output of the combination of $\hat{N}$ atoms with our e-bedding solver and the MP pac solver for the problem of plane wave scattering with three regular polygons— $\hat{N}$ is required $\hat{N}$ and $\hat{N}$ so ves on the triangle square and pentagon respectively a total of $\hat{N}$ so ves for a single scatterer each a number independent of k —For wavenumber $k =$ using MP pac with out solving via the e-bedding for $\hat{N}$ ae results in a total of $\hat{N}$ so ves on each scatterer hence a total of $\hat{N}$ so ves— so even at a relatively low wavenumber the e-bedding for $\hat{N}$ ae can reduce the number of so ves required by $\hat{N}$ atoms as $\hat{N}$ grows with $O(k)$ the number of so ves required by e-bedding for $\hat{N}$ ae is independent of k — there is one added cost with e-bedding for $\hat{N}$ ae namely that the number of terms required in the Taylor series $\hat{N}$ —? and $\hat{N}$ — in order to attain accuracy will need to grow with k we leave detailed consideration of this to future work —

Fig— —, $\hat{N}$ part of total $\hat{N}$ for a configuration of $\hat{N}$ type polygons with incident $\hat{N}$ $u^i_{\pi/}$ so ved using $\hat{N}$ atoms coupled with MP pac and the combined e-bedding approximation of De nition — as the so ver used for the e-bedding presentation w/c in turn is used as the so ver for $\hat{N}$ atoms — the representation is on y valid outside of the union of bases containing each obst

References

- Barnett A–H–and Betc e – ? – tab ty and convergence of t e et od of funda enta so ut ons for He o tz prob e s on ana yt c do a ns–*J. Comput. Phys.* ?? 4 . 4 ?–
- ² Barnett A–H–and Betc e – ? – An exponent a y convergent nonpolyno a n te e e ent et od for t e ar on c scatter ng fro polygons– *SIAM J. Sci. Comput.* ? . 4 44 – MP pac software ava ble to down oad fro github.com/ahbarnett/mpspack–
- B ggs N– ? – A new fa y of e bedd ng for u ae for d ffract on by wedges and polygons–*Wave Motion* 4 ?–
- 4– B ggs N– ? – E bedd ng for u ae for scatter ng